

Some estimates for the viscoelastic incompressible Kelvin-Voigt medium

M.M. Bukenov, D.S. Rakisheva*

*L.N. Gumilyov Eurasian National University, Astana, Kazakhstan
(E-mail: bukenov_m_m@mail.ru, dilya784@mail.ru)*

In this paper, we consider the application of the method of fictitious domains to a viscoelastic incompressible medium based on the Kelvin-Voigt model. Application of the method of fictitious domains allows solving the original problem in regions with complex geometric configuration. This makes it easier to automate the construction of a consistent difference mesh, and to solve the problem in areas of standard shape. Estimates for the proximity of the auxiliary problem's solution are obtained. The auxiliary problem is constructed by the method of fictitious domains. These estimates refer to the solution of the original problem. The original problem describes a viscoelastic incompressible medium. Convergence follows from the estimates of the proximity of the solutions of the original and auxiliary problems. Further, on the basis of the method of fictitious domains, two-sided estimates on a small parameter for the difference between the solution of the original problem and the solution of the auxiliary problem constructed by the method of fictitious domains are obtained. Moreover, the solution to the auxiliary problem is expanded as a series in powers of the small parameter. This is possible because that solution is represented as a functional series that converges absolutely in the original domain.

Keywords: Kelvin-Voigt medium, Kronecker symbol, approximate solution, fictitious area method, two-sided estimates, small parameter, stress, velocity, strain, displacement.

2020 Mathematics Subject Classification: 39A05.

Introduction and problem statement

This paper presents two-sided estimates for a small parameter of the solution of the initial problem. Since a priori estimates help to determine the interval or band in which the solution lies, this is relevant. In addition, since the initial problem does not have an analytical solution, two-sided estimates allow us to determine the initial approximation for finding an approximate solution to the problem, which is an important step in the process of finding a solution.

We consider the application of the method of fictitious domains for incompressible Kelvin-Voigt medium. Two-sided estimates of the convergence of the approximate solution to the exact solution by a small parameter α are obtained. We consider the formulation of a dynamic viscoelastic incompressible medium based on the Kelvin-Voigt model in a cylinder $Q = \{D \times [0 \leq t \leq t_1]\}$, where $D \subset R^3$ is a bounded singly connected region with a sufficiently smooth boundary γ . Let us introduce the notation $\gamma_t = \gamma \times [0, t_1]$, the strain and stress vector-functions $\vec{\varepsilon} = \{\varepsilon_{11}, \varepsilon_{22}, \varepsilon_{33}, 2\varepsilon_{12}, 2\varepsilon_{13}, 2\varepsilon_{23}\}^T$, $\vec{\sigma} = \{\sigma_{11}, \sigma_{22}, \sigma_{33}, \sigma_{12}, \sigma_{13}, \sigma_{23}\}^T$, here the symbol T is transpose, the displacement and velocity vector functions $\vec{U} = \{U_1, U_2, U_3\}^T$, $\vec{v} = \{v_1, v_2, v_3\}^T$. Following the work of [1] we consider the velocity-stress formulation. We find the solution satisfying the following relations:

$$\frac{\partial \vec{v}}{\partial t} + R\vec{\sigma} = \vec{f}, \quad (x, t) \in Q \quad (1)$$

*Corresponding author. E-mail: dilya784@mail.ru

Received: 24 October 2024; Accepted: 28 March 2025.

© 2025 The Authors. This is an open access article under the CC BY-NC-ND license (<http://creativecommons.org/licenses/by-nc-nd/4.0/>)

which is an equation of motion,

$$\frac{\partial \vec{\varepsilon}}{\partial t} - R \vec{\mathcal{J}} = 0, \quad (2)$$

displacement-strain ratio,

$$B \vec{\sigma} = J \vec{\varepsilon} + D \frac{\partial \vec{\varepsilon}}{\partial t}. \quad (3)$$

The equation of state for the Kelvin-Voigt medium is

$$\operatorname{div} \vec{u} = 0. \quad (4)$$

The condition of incompressibility of the medium, taking into account the stresses and the pressure function p , is given by the relationship

$$\sigma_{ik} = -\delta_{ik}p + 2\mu\varepsilon_{ik}, \quad i, k = 1, 2, 3. \quad (5)$$

Here δ_{ik} is the Kronecker symbol, $\vec{f}$ is the vector of mass forces, $B = B^T$, $C = C^T$ are symmetric positive-definite matrices depending on the Lamé constants and viscosity coefficient, J is an diagonal matrix, their form is given in [2].

R is a linear matrix-differential operator:

$$R = \begin{pmatrix} \nabla_1 & 0 & 0 & \nabla_2 & \nabla_3 & 0 \\ 0 & \nabla_2 & 0 & \nabla_1 & 0 & \nabla_3 \\ 0 & 0 & \nabla_3 & 0 & \nabla_1 & \nabla_2 \end{pmatrix}^T, \quad R^* = -R^T, \quad \nabla_i = \frac{\partial}{\partial x_i}, \quad i = 1, 2, 3.$$

The system of equations (1)–(5) transform to the following form, a vector function $\vec{\sigma}(x, t)$ that satisfies the following relations

$$B \frac{\partial^2 \vec{\sigma}}{\partial t^2} = A \vec{\sigma} + DA \frac{\partial \vec{\sigma}}{\partial t} + \vec{F}, \quad (6)$$

$A = -RR^*$, $\vec{F} = R \vec{f}$, satisfies the initial conditions:

$$\vec{\sigma}(x, 0) = \vec{q}(x), \quad \frac{\partial \vec{\sigma}}{\partial t}(x, 0) = \vec{g}(x) \quad (7)$$

and boundary conditions

$$\sum_{k=1}^3 \sigma_{ik}(x, t) n_k = 0, \quad (x, t) \in \gamma_t; \quad (8)$$

here $n_k = (n_1, n_2, n_3)^T$ is the vector of normal to γ , $\gamma_t = \gamma \times [0, t_1]$. Let us denote the problem (6)–(8) by the problem I.

Main provisions

In [2], we show the stationalization of the solution of the problem I to the solution of the static elastic problem.

$$R^* \vec{\sigma}^y(x) + \vec{F}^y(x) = 0, \quad \vec{\sigma}^y(x) = R \varepsilon^y(x),$$

$$\sum_{k=1}^3 \sigma_{ik}^y(x) n_k = 0, \quad x \in \gamma. \quad (9)$$

In [2], the closeness estimate of the solution of problem I and problem (9) is obtained:

$$\|\vec{\sigma} - \vec{\sigma}^y\| \leq e^{-\beta t} \cdot \|\vec{\sigma}(x, 0) - \vec{\sigma}^y(x)\|,$$

where $\beta > 0$ is a constant.

The following theorem is true for problem I:

Theorem 1. Let

$$\vec{\sigma}(x, t) \in W_2^{2,1}(Q), \quad \vec{g}(x) \in \dot{W}_2^1(D), \quad \vec{q}(x) \in L_2(D), \quad \vec{F}(x, t) \in L_2(Q),$$

then there exists a unique solution to the problem I and the following estimation is true

$$\|\vec{\sigma}(x, t)\|_{W_2^{2,1}(Q)} \leq C_1(\|\vec{F}\|_{L_2(Q)} + \|\vec{q}\|_{L_2(D)} + \int_0^{t_1} \|\vec{g}\|_{\dot{W}_2^1(D)} d\tau).$$

The proof is similar to the proof of Theorem 1 in [3].

According to the method of fictitious domains [3–6], we augment the original region D with the region D_1 to a composite region $D_0 = D \cup D_1$, with boundary $\Gamma, \Gamma_t = \Gamma \times [0, t_1]$, $+Q_1 = D_1 \times [0, t_1]$ and construct the auxiliary problem

$$\begin{aligned} L_\alpha \vec{\sigma}^\alpha &= \vec{F}, \quad (x, t) \in Q, \quad L_\alpha \vec{\sigma}^\alpha = 0, \quad (x, t) \in Q_1, \\ \sum_{k=1}^3 (\vec{\sigma}^\alpha)_{ik} n_k &= 0, \quad (x, t) \in \gamma_t, \quad \vec{\sigma}^\alpha(x, 0) = 0, \quad x \in D_1, \\ \vec{\sigma}^\alpha(x, 0) &= \vec{q}(x), \quad x \in D, \quad \frac{\partial \vec{\sigma}^\alpha}{\partial t}(x, 0) = 0, \quad x \in D_1 \\ \left. \frac{\partial \vec{\sigma}^\alpha}{\partial t} \right|_t &= \vec{g}(x), \quad x \in D, \quad \sum_{k=1}^3 \sigma_{ik}^\alpha n_k = 0, \quad (x, t) \in \Gamma_t, \end{aligned} \quad (10)$$

where $L_\alpha \vec{\sigma}^\alpha = B \frac{\partial^2 \vec{\sigma}^\alpha}{\partial t^2} - a^\alpha A \vec{\sigma}^\alpha + JA \frac{\partial \vec{\sigma}^\alpha}{\partial t}$, $a^\alpha = \begin{cases} 1, & x \in D, \\ -\alpha^2, & x \in D_1, \end{cases}$ $\alpha > 0$ is a small parameter.

On the coefficient gap curve γ_t , we set the following matching conditions

$$\vec{\sigma}^\alpha|_{\gamma_t}^+ = \vec{\sigma}^\alpha|_{\gamma_t}^-, \quad \left. \frac{\partial \vec{\sigma}^\alpha}{\partial N} \right|_{\gamma_t}^+ = \frac{M}{\alpha} \left. \frac{\partial \vec{\sigma}^\alpha}{\partial n} \right|_{\gamma_t}^-.$$

Signs “+” or “−” mean convergence to the limit value of the function from inside or outside to the boundary γ_t . The parameter M takes the values -1 or $+1$ [6, 7].

Let us introduce the following series into consideration:

$$S_1 = \sum_{k=0}^{\infty} \alpha^k \vec{V}_k, \quad \text{on } Q, \quad S_2 = \sum_{k=1}^{\infty} \alpha^k \vec{W}_k, \quad \text{on } Q_1. \quad (11)$$

Putting (11) into (10), we obtain relations for determining $\vec{V}_k$ and $\vec{W}_k$:

$$\begin{aligned} L_\alpha \vec{V}_0 &= \vec{F}, \quad (x, t) \in Q, \quad \vec{W}_1 = 0, \quad (x, t) \in Q_1, \\ \sum_{k=1}^3 (V_0)_{ik} n_k &= 0, \quad (x, t) \in \gamma_t, \quad \vec{W}_1(x, 0) = 0, \quad \frac{\partial \vec{W}_1}{\partial t}(x, 0) = 0, \quad x \in D_1. \\ \vec{V}_0(x, 0) &= \vec{q}(x), \quad \frac{\partial \vec{V}_0}{\partial t}(x, 0) = \vec{g}(x), \quad x \in D, \quad \frac{\partial \vec{W}_1}{\partial n} = M \frac{\partial \vec{V}_0}{\partial N}, \quad (x, t) \in \gamma_t. \end{aligned} \quad (12)$$

$$\sum_{k=1}^3 (\vec{W}_1)_{ik} n_k = 0, \quad (x, t) \in \gamma_t.$$

And for $k \geq 1$

$$L_\alpha \vec{V}_k = \vec{F}, \quad (x, t) \in Q, \quad L_\alpha \vec{W}_{k+1} = 0, \quad (x, t) \in Q_1,$$

$$\sum_{k=1}^3 (V_k)_{ik} n_k = 0, \quad (x, t) \in \gamma_t, \quad \sum_{k=1}^3 (W_{k+1})_{ik} n_k = 0, \quad (x, t) \in \gamma_t,$$

$$\vec{V}_k(x, 0) = 0, \quad x \in D, \quad \vec{W}_{k+1}(x, 0) = 0, \quad x \in D_1,$$

$$\frac{\partial \vec{V}_k}{\partial t}(x, 0) = 0, \quad x \in D, \quad \frac{\partial \vec{W}_{k+1}}{\partial t}(x, 0) = 0, \quad x \in D_1,$$

$$\vec{V}_k = \vec{W}_k, \quad (x, t) \in \gamma_t.$$

Functions $\vec{V}_k \in W_2^{2,1}(Q)$, $k = 0, 1, \dots$, $\vec{W}_k \in W_2^{2,1}(Q_1)$, $k = 0, 1, \dots$

We obtain estimates of the convergence of the solution of the auxiliary problem to the solution of the original problem with respect to the small parameter α .

Theorem 2. If α_0 is such that $0 < \alpha < \alpha_0$, then the series S_1, S_2 absolutely converge to $W_2^{2,1}(Q)$ and $W_2^{2,1}(Q_1)$, and the following equalities are true

$$\vec{\sigma}^\alpha = S_1, \quad (x, t) \in Q, \quad \vec{\sigma}^\alpha = S_2, \quad (x, t) \in Q_1,$$

where $\vec{\sigma}^\alpha$ is the solution of problem (10).

Proof. We search obvious priori estimates [7, 8].

$$\|\vec{W}_k\|_{W_2^{2,1}(Q_1)} \leq C_1 \left\| \frac{\partial \vec{W}_k}{\partial n} \right\|_{W_2^{\frac{1}{2}}(\gamma_t)} \leq C_2 \left\| \frac{\partial \vec{V}_{k-1}}{\partial N} \right\|_{W_2^{\frac{1}{2},1}(\gamma_t)} \leq C_1 C_2 \|\vec{V}_{k-1}\|_{W_2^{\frac{1}{2}}(\gamma_t)}, \quad (13)$$

where C_1, C_2 are constants depending on the regions D, D_1 and not depending on α .

Now we show the convergence of the series S_1 to $W_2^{2,1}(Q_1)$ and S_2 to $W_2^{2,1}(Q_1)$, we have

$$\|\vec{V}_k\|_{W_2^{2,1}(Q)} \leq C_3 \|\vec{V}_k\|_{W_2^{\frac{3}{2},1}(\gamma_t)} = C_3 \|\vec{W}_k\|_{W_2^{\frac{3}{2},1}(\gamma_t)} \leq C_3 C_4 \|\vec{W}_k\|_{W_2^{2,1}(Q_1)},$$

and using (8), (13), we obtain

$$\|\vec{V}_k\|_{W_2^{2,1}(Q)} \leq C_5 \|\vec{V}_{k-1}\|_{W_2^{\frac{3}{2},1}(Q)}, \quad k \geq 1,$$

then

$$\|\vec{V}_0\|_{W_2^{2,1}(Q)} \leq C_6 \left(\|\vec{F}\|_{L_2(Q)} + \|\vec{q}\|_{L_2(D)} + \int_0^{t_1} \|\vec{g}\|_{\dot{W}_2^1(D)} dt \right),$$

where $C_6 = C_1 C_2 C_3 C_4 C_5$.

Assuming $\alpha < \alpha_0 = C_6^{-1}$, we obtain the series S_1 , is absolutely convergent to $W_2^{2,1}(Q)$ and correspondingly the series S_2 is absolutely convergent to $W_2^{2,1}(Q_1)$.

Multiplying (12) for $\vec{V}_k$ and $\vec{W}_k$ by α_k , and summing over k , we have

$$\begin{aligned} LS_1 &= \vec{F}, \quad (x, t) \in Q, \quad L_\alpha S_2 = 0, \quad (x, t) \in Q_1, \\ S_1(x, 0) &= \vec{q}(x), \quad x \in D, \quad S_2(x, 0) = 0, \quad x \in D_1, \end{aligned} \quad (14)$$

$$\begin{aligned}\frac{\partial S_1}{\partial t}(x, 0) &= \vec{g}(x), \quad x \in D, \quad \frac{\partial S_2}{\partial t}(x, 0) = 0, \quad x \in D_1. \\ S_2(x, t) &= 0, \quad (x, t) \in \gamma_t, \\ S_1 &= S_2, \quad (x, t) \in \Gamma_t, \quad \frac{\partial S_2}{\partial n} = \frac{M}{\alpha} \frac{\partial S_1}{\partial N}, \quad (x, t) \in \gamma_t, \\ L\vec{\sigma} &= B \frac{\partial^2 \vec{\sigma}}{\partial t^2} - A\vec{\sigma} - JA \frac{\partial \vec{\sigma}}{\partial t}.\end{aligned}$$

Hence, (14) we obtain that $\vec{\sigma}^\alpha = S_1$ in Q and $\vec{\sigma}^\alpha = S_2$ in Q_1 , if the condition $0 < \alpha < \alpha_0$ is met. From the proof of Theorem 2 implies the following statement

$$\|\vec{\sigma} - \vec{\sigma}_+^\alpha\|_{W_2^{2,1}(Q)} \leq C_7\alpha, \quad \|\vec{\sigma} - \vec{\sigma}_-^\alpha\|_{W_2^{2,1}(Q)} \leq C_8\alpha, \quad (15)$$

here $\vec{\sigma}_+^\alpha = \vec{\sigma}^\alpha$, at $M = 1$, $\vec{\sigma}_-^\alpha = \vec{\sigma}^\alpha$, at $M = -1$. $\vec{\sigma}$ is the solution to the problem I. C_7, C_8 are constants depend on the areas D, D_1 and are independent of α .

Next, we can formulate a theorem giving two-sided estimates on α [9].

Theorem 3. If $0 < \alpha < \alpha_0$, $\vec{\sigma}$ is a solution of problem I, $\vec{\sigma}_+^\alpha$, $\vec{\sigma}_-^\alpha$ is a solution of problem (10) at $M = 1$ and $M = -1$, then the following estimation (16) is true

$$\left\| \vec{\sigma} - \frac{1}{2} (\vec{\sigma}_+^\alpha + \vec{\sigma}_-^\alpha) \right\|_{W_2^{2,1}(Q)} \leq C_9\alpha^2, \quad (16)$$

where

$$\vec{\sigma}^\alpha = S_1, \quad (x, t) \in Q, \quad \vec{\sigma}^\alpha = S_2, \quad (x, t) \in Q_1. \quad (17)$$

Proof. By virtue of Theorem 2, we have

$$\vec{\sigma}_+^\alpha = \sum_{k=0}^{\infty} \alpha^k \vec{V}_k^+, \quad (x, t) \in Q, \quad \vec{\sigma}_+^\alpha = \sum_{k=1}^{\infty} \alpha^k \vec{W}_k^+, \quad (x, t) \in Q_1, \quad (18)$$

here $\vec{V}_k^+, \vec{W}_k^+$ are solutions of (10) at $M = 1$, moreover

$$\vec{\sigma}_-^\alpha = \sum_{k=0}^{\infty} \alpha^k \vec{V}_k^-, \quad (x, t) \in Q, \quad \vec{\sigma}_-^\alpha = \sum_{k=1}^{\infty} \alpha^k \vec{W}_k^-, \quad (x, t) \in Q_1, \quad (19)$$

here $\vec{V}_k^-, \vec{W}_k^-$ are solutions of (10) at $M = -1$.

We obtain $\vec{V}_0^+ \equiv \vec{V}_0^- \equiv \vec{\sigma}$, it is a solution of problem I.

We introduce the notation $\vec{W}_1 = \vec{W}_1^+ + \vec{W}_1^-$, the function $\vec{W}_1$ satisfies the following problem

$$\begin{aligned}L_\alpha \vec{W}_1 &= 0, \quad (x, t) \in Q_1, \quad \frac{\partial \vec{W}_1}{\partial n} = 0, \quad (x, t) \in \gamma_t, \\ \vec{W}_1(x, 0) &= 0, \quad \frac{\partial \vec{W}_1}{\partial t}(x, 0) = 0, \quad x \in D_1, \quad \vec{W}_1(x, t) = 0, \quad (x, t) \in \Gamma_t,\end{aligned}$$

hence, we obtain that $\vec{W}_1 = 0$, or $\vec{W}_1^+ = -\vec{W}_1^-$.

Further we introduce $\vec{V}_1 = \vec{V}_1^+ + \vec{V}_1^-$, the function $\vec{V}_1$ satisfies the problem

$$L_\alpha \vec{V}_1 = 0, \quad (x, t) \in Q, \quad \frac{\partial \vec{V}_1}{\partial t}(x, 0) = 0, \quad x \in D,$$

$$\vec{V}_1(x, 0) = 0, \quad (x, t) \in \Gamma_t, \quad \vec{V}_1(x, t) = 0, \quad (x, t) \in \gamma_t,$$

from which we obtain $\vec{V}_1 = 0$, or $\vec{V}_1^+ = -\vec{V}_1^-$. By sequentially introducing

$$\vec{W}_2 = \vec{W}_2^+ + \vec{W}_2^-, \quad \vec{V}_2 = \vec{V}_2^+ + \vec{V}_2^-,$$

we get

$$\vec{W}_2 = \vec{W}_2^-, \quad \vec{V}_2^+ = \vec{V}_2^-,$$

continuing this process, we get

$$\begin{aligned} \vec{V}_k^+ &= \vec{V}_k^-, \text{ if } k \text{ is even,} \\ \vec{V}_k^+ &= -\vec{V}_k^-, \text{ if } k \text{ is odd.} \end{aligned}$$

Substituting (20) into (18), (19), we have

$$\vec{\sigma}_+^\alpha = \vec{\sigma} + \sigma \vec{V}_1^+ + \sigma^2 \vec{V}_2^+ + \dots$$

$$\vec{\sigma}_-^\alpha = \vec{\sigma} - \sigma \vec{V}_1^+ + \sigma^2 \vec{V}_2^+ + \dots$$

Applying decomposition (20), as well as estimation (17) at $0 < \alpha < \alpha_0$, we obtain

$$\left\| \vec{\sigma} - \frac{1}{2} (\vec{\sigma}_+^\alpha + \vec{\sigma}_-^\alpha) \right\|_{W_2^{2,1}(Q)} \leq \alpha^2 \left\| \vec{V}_2^+ + \alpha^2 \vec{V}_4^+ + \dots \right\|_{W_2^{2,1}(Q)} \leq C_8 \alpha^2 \left\| \vec{V}_0^+ \right\|_{W_2^{2,1}(Q)} \leq C_9 \alpha^2;$$

here $C_8 = C_\alpha^2$, so for $x \in D$, $0 < \alpha < \alpha_0$, we have a two-sided estimation

$$O(\alpha^2) + \min(\vec{\sigma}_+^\alpha, \vec{\sigma}_-^\alpha) \leq \vec{\sigma} \leq \max(\vec{\sigma}_+^\alpha, \vec{\sigma}_-^\alpha) + O(\alpha^2).$$

Thus, a two-sided estimate in terms of the small parameter of the solution to the original problem has been obtained through the solution of the auxiliary problem, where the parameter values $M = -1$, $M = 1$ corresponds to $\vec{\sigma}_+^\alpha$.

Conclusion

The obtained estimate is essential for the application of the numerical solution of the auxiliary problem, and it is not considered in the works [10–16]. Continuation by the lowest coefficient in the fictitious region method leads to the same estimates. In works [10–16], the numerical implementation of the Kelvin-Voigt model in equivalent formulations is considered.

Author Contributions

All authors contributed equally to this work.

Conflict of Interest

The authors declare no conflict of interest.

References

- 1 Bukenov, M.M. (2005). Postanovka dinamicheskoi zadachi lineinoi viazkouprugosti v skorostiakh-napriazheniiakh [Formulation of a dynamic problem of linear viscoelasticity in terms of velocities and stresses]. *Sibirskii zhurnal vychislitelnoi matematiki — Siberian Journal of Computational Mathematics*, 8(4), 289–295 [in Russian].
- 2 Patsyuk, V.I. (1982). Statsionirovanie dinamicheskikh protsessov v viazkouprugikh sredakh [Stationarization of dynamic processes in viscoelastic media]. *Candidate's thesis* [in Russian].
- 3 Bukenov, M.M., & Azimova, D.N. (2019). Estimates for Maxwell viscoelastic medium in tension-rates. *Eurasian Mathematical Journal*, 10(2), 30–36. <http://dx.doi.org/10.32523/2077-9879-2019-10-2-30-36>
- 4 Konovalov, A.N. (1972). Metod fiktivnykh oblastei v zadachakh filtratsii dvukhfaznoi neszhimaemoi zhidkosti s uchetom kapilliarnykh sil [The method of fictitious domains in problems of two-phase incompressible fluid filtration considering capillary forces]. *Chislennye metody mekhaniki sploshnoi sredy — Numerical Methods of Continuum Mechanics*, 3(5), 52–67 [in Russian].
- 5 Konovalov, A.N. (1975). Ob odnom variante metoda fiktivnykh oblastei [On one version of the fictitious region method]. *Nekotorye problemy vychislitelnoi i prikladnoi matematiki — In Some problems of computational and applied mathematics* (pp. 191–199). Novosibirsk [in Russian].
- 6 Orunkhanov, M.K., & Smagulov, Sh. (1982). K teorii metoda fiktivnykh oblastei [To the theory of the fictitious domain method]. *Chislennye metody mekhaniki sploshnoi sredy — Numerical Methods of Continuum Mechanics*, 13(2), 125–137 [in Russian].
- 7 Ladyzhenskaya, O.A. (1977). *Kraevye zadachi matematicheskoi fiziki [Boundary Problems of Mathematical Physics]*. Moscow: Nauka [in Russian].
- 8 Lions, J.L., & Magenes, E. (1972). *Non-Homogeneous Boundary Value Problems and Applications, I*. Springer-Verlag, Berlin, Heidelberg, New York.
- 9 Chertova, K. (1994). Locally two-sided approximate solutions in parabolic problems. *Bulletin of the Novosibirsk Computing Center. Series: Numerical analysis*, 6, 37–42.
- 10 Zhang, R., & He, Q. (2022). The least-square/fictitious domain method based on Navier slip boundary condition for simulation of flow–particle interaction. *Applied Mathematics and Computation*, 415(1), 126687. <https://doi.org/10.1016/j.amc.2021.126687>
- 11 Mukanova, B., Turarova, M., & Rakisheva, D. (2025). Longitudinal electrical resistivity tomography interpretation: Numerical modelling. *Geophysical Prospecting*, 73(2), 680–698. <https://doi.org/10.1111/1365-2478.13614>
- 12 Bazilevs, Y., Takizawa, K., & Tezduyar, T.E. (2013). *Computational Fluid-Structure Interaction: Methods and Applications*. New York: Wiley. <http://dx.doi.org/10.1002/9781118483565>
- 13 Wu, M., Peters, B., Rosemann T., & Kruggel–Emden, H. (2020). A forcing fictitious domain method to simulate fluid–particle interaction of particles with super–quadric shape. *Powder Technology*, 360(1), 264–277. <https://doi.org/10.1016/j.powtec.2019.09.088>
- 14 Boffi, D., Gong, S., Guzmán, J., & Neilan, M. (2024). Convergence of Lagrange finite element methods for Maxwell eigenvalue problem in 3D. *IMA Journal of Numerical Analysis*, 44(4), 1911–1945. <https://doi.org/10.1093/imanum/drad053>
- 15 Nobile, F., & Vergara, C. (2008). An effective fluid-structure interaction formulation for vascular dynamics by generalized Robin conditions. *SIAM Journal on Scientific Computing*, 30(2), 731–763. <https://doi.org/10.1137/060678439>
- 16 Rakisheva, D.S., & Mukanova, B.G. (2021). Fourier Transformation method for solving integral equation in the 2.5D problem of electric sounding. *Journal of Physics: Conference Series*, 2092, 012018. <https://doi.org/10.1088/1742-6596/2092/1/012018>

*Author Information**

Makhat Muchamedyevich Bukenov — Associate Professor of Physical and Mathematical Sciences, Department of Mathematical and Computer Modelling, Faculty of Mechanics and Mathematics, L.N. Gumilyov Eurasian National University, Astana, 010008, Kazakhstan; e-mail: bukenov_m_m@mail.ru; <https://orcid.org/0000-0003-4813-8677>

Dilyara Sovetovna Rakisheva (*corresponding author*) — Head of the Department of Mathematical and Computer Modeling Faculty of Mechanics and Mathematics, L.N. Gumilyov Eurasian National University, Astana, 010008, Kazakhstan; e-mail: dilya784@mail.ru; <https://orcid.org/0000-0001-8434-9469>

*The author's name is presented in the order: First, Middle, and Last Names.